

Zeroth order moment

$$\int \frac{\partial f}{\partial t} d^3\bar{v} + \int \bar{v} \cdot \frac{\partial f}{\partial \bar{x}} d^3\bar{v} + \frac{q}{m} \int (\bar{E} + \bar{v} \times \bar{B}) \frac{\partial f}{\partial \bar{v}} d^3\bar{v} = \int \left(\frac{\partial f}{\partial t} \right)_c d^3\bar{v}$$

First term

$$\int \frac{\partial f}{\partial t} d^3\bar{v} = \frac{\partial}{\partial t} \int f d^3\bar{v} = \frac{\partial n}{\partial t}$$

$$\int \bar{v} \cdot \frac{\partial f}{\partial \bar{x}} d^3\bar{v} = \frac{1}{\partial \bar{x}} \int \bar{v} d^3\bar{v} = \frac{\partial}{\partial \bar{x}} (n \bar{v}) = \frac{\partial}{\partial \bar{x}} n \bar{v}$$

$$\int (\bar{E} \cdot \frac{\partial f}{\partial \bar{v}}) d^3\bar{v} = \int \frac{\partial}{\partial \bar{v}} (f \bar{E}) d^3\bar{v} = \int_S f \bar{E} \cdot d\bar{S}_v = 0$$

Gauss' law

Vanishes as we take $S \rightarrow \infty$

$S \xrightarrow{\text{incre}} v^2$ but $f(\bar{v})$ med distributed
will go to zero $\exp(-v^2/v_T^2)$

$$\int (\bar{v} \times \bar{B}) \cdot \frac{\partial f}{\partial \bar{v}} d^3\bar{v} = \int \frac{\partial}{\partial \bar{v}} (f (\bar{v} \times \bar{B})) d^3\bar{v}$$

$$= \int f \frac{\partial}{\partial \bar{v}} \cdot (\bar{v} \times \bar{B}) d^3\bar{v} = 0$$

First term as surface integral

Second term $\bar{v} \times \bar{B} \perp \bar{v}$

Collision term also goes to zero since number of particles of specie remain constant.

$$\frac{\partial n}{\partial t} + \frac{\partial}{\partial x} (n \bar{v}) = 0$$

Conservation of particles.

Second order. multiply by $m \bar{v}$

$$m \int \bar{v} \frac{\partial f}{\partial t} d^3 \bar{v} + m \int \bar{v} \left(\bar{v} \cdot \frac{\partial}{\partial x} \right) f d^3 \bar{v} \\ + q \int \bar{v} (\bar{E} + \bar{v} \times \bar{B}) \frac{\partial f}{\partial \bar{v}} d^3 \bar{v} \\ = \int m \bar{v} \left(\frac{\partial f}{\partial t} \right) d^3 \bar{v}$$

The first term gives

$$m \int \bar{v} \frac{\partial f}{\partial t} d^3 \bar{v} = m \frac{\partial}{\partial t} \int \bar{v} d^3 \bar{v} = m \frac{\partial n \bar{v}}{\partial t}$$

$$- q \int (\bar{E} + \bar{v} \times \bar{B}) f d^3 \bar{v} = - q n (\bar{E} + \bar{v} \times \bar{B})$$

Second term.

$$\frac{\partial}{\partial x} \cdot (n \bar{v} \bar{v})$$

$$\mu = \left(\frac{F}{A} \right) / \frac{\Delta U_x}{\Delta z}$$

①

from momentum eqⁿ

assume const density & pressure
& ignore any external forces.

$$\frac{\partial}{\partial t} (\rho U_x) = \frac{\partial}{\partial y} \underbrace{\mu \frac{\partial U_x}{\partial y}}$$

Viscous force per unit volume acting on
fluid element

The rate at which particles cross
upper surface

$$\sim \frac{n v_T}{2} \Delta A \quad v_T \sim \left(\frac{kT}{m} \right)^{1/2}$$

The particles travel distance λ before colliding
& exchanging momentum, so even though
equal number of particles cross the boundary
in upward & downward. there is
systematic difference in momentum

$$dp_x \sim 2m\lambda \frac{\partial U_x}{\partial y} \quad \text{per particle.}$$

The rate of change of the element momentum
at boundaries.

$$\frac{dp_x}{dt} \sim \left[m\lambda \frac{\partial U_x}{\partial y} n v_T \Delta A \right]_{@ y+\Delta y} - \left[m\lambda \frac{\partial U_x}{\partial y} n v_T \Delta A \right]_{@ y}$$

we must subtract momentum taken out
by particles crossing the bottom surface
Dividing $DA dy$ by the volume. (2)

force per unit volume

$$\frac{\partial}{\partial y} \left(\mu \frac{\partial u_x}{\partial y} \right) \sim \frac{\partial}{\partial y} \left(n m \lambda v_T \frac{\partial u_x}{\partial y} \right)$$

using

$$\lambda = \frac{1}{n \sigma}$$

$$\mu \approx \frac{m v_T}{\sigma}$$

$$\eta = \frac{1}{3\sqrt{2}} \frac{m}{\sigma} \bar{v}$$

$$\bar{v} = \sqrt{\frac{8}{\pi} \frac{k_B T}{m}}$$

η = depends only on T not on
the density or pressure.

Heat Conduction

$$\vec{Q} = -K \vec{\nabla} T$$

$$Q_z = -k \frac{\partial T}{\partial z}$$

$$K = \frac{1}{3} n \bar{v} c \lambda$$

$c = \frac{\partial \bar{E}}{\partial T}$ Specific energy per molecule

$$K = \frac{1}{3\sqrt{2}} \frac{c}{\sigma_0} \bar{v} \quad \text{independent of } P$$

$$\frac{k}{\eta} = \frac{c}{m}$$

Diffusion

due to density gradient it arise

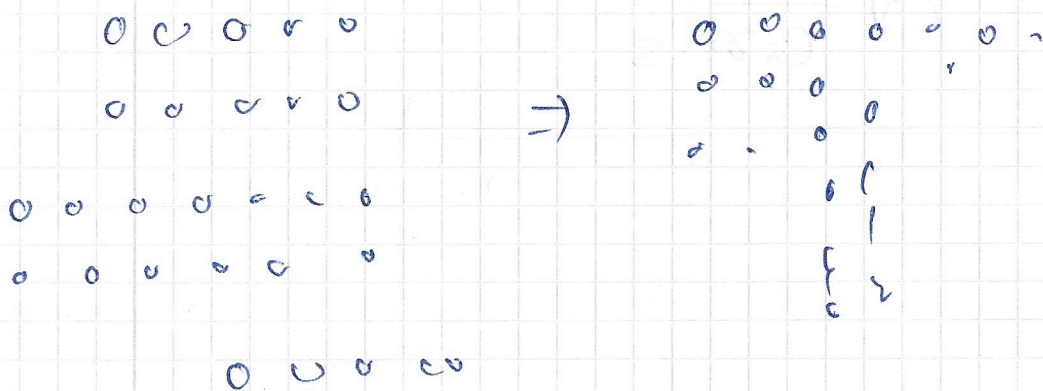
$$j_z = -D \frac{\partial n}{\partial z}$$

$$D = \frac{1}{3} \bar{v} \lambda$$

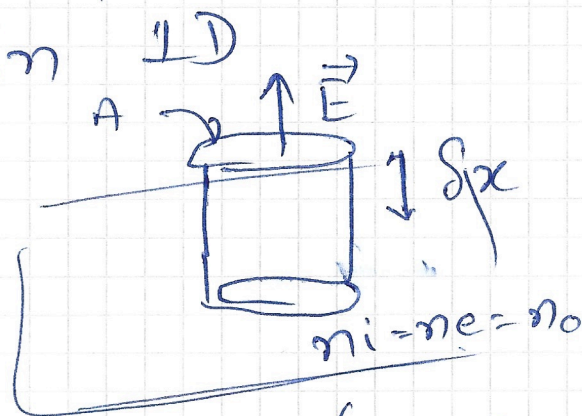
$$\frac{D}{\eta} = \frac{1}{nm}$$

Plasma Oscillations

$$n_{0i} \approx n_{0e} = n_0$$



Consider a small perturbation of electric charge in 1D



$$Q = -e n_0 (A \delta x)$$

$$\int \vec{E} \cdot d\vec{s} = A E_{1x} = \frac{Q}{\epsilon_0} \Rightarrow E_{1x} = -\frac{e}{\epsilon_0} n_0 A \delta x$$

$$E_{1x} = -\frac{e}{\epsilon_0} n_0 \delta x$$

given eqn of motion

$$\frac{d^2(\delta x)}{dt^2} + \frac{e^2 n_0}{m_s \epsilon_0} (\delta x) = 0$$